

**BE/APh 161: Physical Biology of the Cell, Winter 2019**  
**Homework #7**

Due at the start of lecture, 2:30 PM, February 27, 2019.

**Problem 7.1** (Genomes in cells, 10 pts).

In this problem we consider how genomes take up space in cells.

- a) If your genome were a single strand of DNA, what would its approximate radius of gyration be if it were free in solution? What implications does this have for the design of a cell?
- b) Estimate the radius of gyration of the *E. coli* genome if it were not confined in a cell. How does that compare to the size of an *E. coli* cell?

**Problem 7.2** (The WLC in the stiff limit, 5 pts).

In class, we worked on deriving the mean squared end-to-end distance of a wormlike chain. I made a sign error in lecture, which resulted in us getting the right limit for with  $\xi_p \ll L$  limit (where  $\xi_p$  is the persistence length and  $L$  is the total length of the polymer), but the wrong limit for  $\xi_p \gg L$ . This is an example where taking limits can be useful in checking your work.

Here, I repeat the derivation of the mean square end-to-end distance to get the correct expression at the end. Remember that  $\mathbf{u}(s)$  is the unit tangent vector to the polymer, and we define the persistence length  $\xi_p$  according to

$$\langle \mathbf{u}(s) \cdot \mathbf{u}(s') \rangle = e^{-|s-s'|/\xi_p}. \quad (7.1)$$

(We later worked out a relationship between  $\xi_p$  and the flexural rigidity,  $K$ ,  $\xi_p = K/k_B T$ .) Now, we compute the mean square end-to-end distance, using the fact that the end-to-end vector is

$$\mathbf{R} = \int_0^L ds \mathbf{u}(s). \quad (7.2)$$

We compute directly.

$$\begin{aligned} \langle \mathbf{R} \cdot \mathbf{R} \rangle &= \left\langle \left( \int_0^L ds \mathbf{u}(s) \right) \cdot \left( \int_0^L ds \mathbf{u}(s) \right) \right\rangle = \left\langle \int_0^L ds \int_0^L ds' \mathbf{u}(s) \cdot \mathbf{u}(s') \right\rangle \\ &= \int_0^L ds \int_0^L ds' \langle \mathbf{u}(s) \mathbf{u}(s') \rangle = \int_0^L ds \int_0^L ds' e^{-|s-s'|/\xi_p} \\ &= 2 \int_0^L ds \int_0^s ds' e^{-(s-s')/\xi_p} = 2 \int_0^L ds e^{-s/\xi_p} \int_0^s ds' e^{s'/\xi_p} \\ &= -2 \xi_p \int_0^L ds e^{-s/\xi_p} (1 - e^{s/\xi_p}) = 2 \xi_p \int_0^L ds (1 - e^{-s/\xi_p}) \end{aligned}$$

$$= 2L\xi_p \left( 1 - \frac{\xi_p}{L} (1 - e^{-L/\xi_p}) \right). \quad (7.3)$$

This is the correct expression, and in lecture I erroneously wrote

$$\langle \mathbf{R} \cdot \mathbf{R} \rangle = 2L\xi_p \left( 1 + \frac{\xi_p}{L} (1 - e^{-L/\xi_p}) \right) \quad (\text{WRONG}). \quad (7.4)$$

In lecture, we showed (correctly, but with the incorrect full expression for  $\langle \mathbf{R} \cdot \mathbf{R} \rangle$ ) that in the flexible chain limit,  $L \gg \xi_p$ , that  $\langle \mathbf{R} \cdot \mathbf{R} \rangle \approx 2L\xi_p$ . Now, show that in the stiff limit ( $L \ll \xi_p$ ),

$$\langle \mathbf{R} \cdot \mathbf{R} \rangle = L^2 \left( 1 - \mathcal{O} \left( \frac{L}{\xi_p} \right) \right). \quad (7.5)$$

**Problem 7.3** (Viral packaging, 20 pts).

In this problem, we explore estimates of the energetics of viral packaging of  $\phi 29$ , which we introduced in lecture.

- a) (Based on problem 10.6 of *PBoC2*) Estimate the entropy penalty for packing the genome in the viral capsid. You can assume that the entropy of the packed state is nearly zero, since it features almost crystalline packaging. Compare this entropy contribution to the free energy (equal to the entropy times the temperature) to the total bending free energy of packing given by equation 10.42 of *PBoC2*. *Hint:* In computing the entropy of the unpacked state, remember that each configuration of the unpacked state has the same energy, since it is a flexible chain on the length scale of the entire genome.
- b) Compare the force required to pack the last bits of genome into the capsid as given by equation 10.43 of *PBoC2* and by the experimental result of  $\approx 50$  pN discussed in lecture and in Figure 10.19(B) of *PBoC2*. What factors might account for any discrepancy you may notice? Make sure you know how equation 10.43 of *PBoC2* is derived.

**Problem 7.4** (Flexural rigidity of biopolymers, adapted from problem 10.2 of *PBoC2*, 30 points).

- a) Recall that the flexural rigidity of a filament is  $K_{\text{eff}} = EI$ , where  $E$  is the Young's modulus and  $I$  is the geometric moment of inertia defined in lecture. We also saw that the persistence length is given by  $\xi_p = EI/k_B T$ . Given the persistence lengths of DNA, actin filaments, and microtubules (check your lecture notes or BioNumbers), estimate their respective Young's moduli by computing the moment of inertia. You can look up geometric information about the filaments in *PBoC2* sections 2.2.3 and 10.5.1.

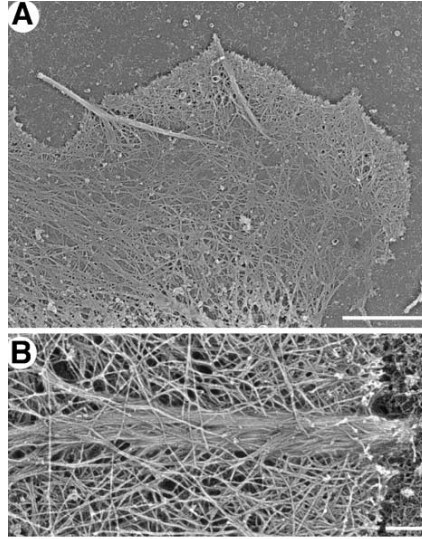


Figure 1: A) Electron micrograph of a B16F1 cell with a few peripherally located filopodia. Scale bar, 5  $\mu\text{m}$ . B) A close-up of one of the filopodia. Scale bar, 1  $\mu\text{m}$ . Image taken from Mejillano, et al., *Cell*, **118**, 363–373, 2004.

- b) Filopodia are protrusions of bundled actin filaments often found in adherent cells. They push against the cell membrane. The membrane pushes back on the filopodium with a force of

$$F = 2\pi r\gamma, \quad (7.6)$$

where  $r$  is the radius of the end of the filopodium and  $\gamma$  is the surface tension of the membrane. Unfortunately, we do not have time this term to talk about membrane tensions, but for this problem, we will take  $\gamma \approx 0.035$  pN/nm. We will assume that the filopodium consists of approximately 30 filaments. We will now investigate how long the filopodium can protrude before it buckles, considering two limits.

- i) First, we assume that the filaments in the filopodium are not crosslinked. Find the length  $L$  that the filopodium can protrude before buckling.
- ii) Now, consider the limit where the filaments in the filopodium are very tightly crosslinked, so tightly crosslinked that the filopodium can be considered a solid rod. Find the length  $L_{\text{cl}}$  that the crosslinked filopodium can protrude before buckling.
- iii) In general, what is  $L_{\text{cl}}/L$  as a function of  $N$ , the number of filaments in the filopodium?

**Problem 7.5** (Polymerization as a force generator, 35 pts).

Imagine an actin filament is polymerizing against a compressive force. This might

be the case if it polymerizes against a membrane, which can deform but nonetheless provides a compressive force on the filament.

- a) Let  $K_d$  be the dissociation constant for binding an additional actin monomer to the end of an actin filament, as defined in lecture. Let  $\delta$  be the increased length of an actin filament as a result of adding one monomer. Show that at equilibrium, the filament can exert a force of

$$F_{\text{eq}} = \frac{k_B T}{\delta} \ln \frac{c_1}{K_d}, \quad (7.7)$$

where  $c_1$  is the concentration of actin monomer. Estimate  $F_{\text{eq}}$  for actin, given that cells typically have  $c_1 \approx 20 \mu\text{M}$ . *Hint:* It might help to think about states and weights.

- b) What is the maximal length of a filament such that it can polymerize against a compressive load without buckling? Derive an analytical expression and then plug in numbers for actin.
- c)  $F_{\text{eq}}$  is the maximal force a filament can exert against a compressive load, as at equilibrium the polymerization force balances the compressive load. Experimentally, it is often the case that this force is never achieved, with polymerization essentially stalling at forces smaller than  $F_{\text{eq}}$ . Provide an intuitive explanation as to why this might be the case.